## Research Article

# APPLICATION OF MATRIX IN STUDY OF NAVIER-STOKES EQUATIONS Patil D. R* 

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## Key Words:

Navier-stokes equation, matrix form, Hassian matrix.


#### Abstract

In present paper we are transforming well known Navier-Stokes equations from linear system to it's matrix form. We are discussing some of the basic properties using properties of matrices such as determinant, eigen value, trace. Using Hassian matrix, we are finding critical points of motion of fluid.


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## INTRODUCTION

The well-known equations of motion of fluid in $\mathrm{R}^{3}$ are the Naiver-Stokes equations. They are second ordered non-linear partial differential equations. Let the Euclidean space be coordinatising with space co-ordinates ( $\mathrm{x}, \mathrm{y}, \mathrm{z}$ ). Assume that fluid is moving in $\mathrm{R}^{3}$ has a fluid element S with position vector ( $\mathrm{x}, \mathrm{y}$, $z$ ) at given specific time. Let $U=\left(u_{x}, u_{y}, u_{z}\right)$ represents the velocity of particle $S$ and $p$ denotes pressure exerted by fluid on $S$ due to it's motion. Let $\rho$ denote the density of fluid, Re denotes is Reynolds number, $E_{t}$ denotes total energy , $\mathrm{P}_{\mathrm{r}}$ denotes Prandtl number and $\tau$ denote amount of flux transport through unit area having nine components along all three directions. Here $\mathrm{x}, \mathrm{y}, \mathrm{z}$ and t are free variables and six unknown quantities as $\mathrm{P}=$ pressure, $\rho=$ density, $\mathrm{T}=$ temperature, $\mathrm{u}_{\mathrm{x}}=$ velocity component along x direction, $\mathrm{u}_{\mathrm{y}}=$ velocity component along y direction, $\mathrm{u}_{\mathrm{z}}=$ velocity component along z direction. So due to conservation of mass, momentum and energy and using Newton's II law we can write Navier-Stokes Equations in $\mathrm{R}^{3}$ as $\frac{\partial\left(\rho u_{x}\right)}{\partial t}+\frac{\partial\left(\rho u_{x}^{2}\right)}{\partial x}+\frac{\partial\left(\rho u_{x} u_{y}\right)}{\partial y}+\frac{\partial\left(\rho u_{x} u_{z}\right)}{\partial z}=-\frac{\partial p}{\partial t}+\frac{1}{R e_{r}}\left[\frac{\partial \tau_{x x}}{\partial x}+\frac{\partial \tau_{x y}}{\partial y}\right.$
$\left.+\frac{\partial \tau_{x z}}{\partial z}\right] \quad$ (along x-axis)
$\frac{\partial\left(\rho u_{y}\right)}{\partial t}+\frac{\left.\partial\left(\rho u_{x} u_{y}\right)\right)}{\partial x}+\frac{\partial\left(\rho u_{y}^{2}\right)}{\partial y}+\frac{\partial\left(\rho u_{y} u_{z}\right)}{\partial z}=-\frac{\partial p}{\partial t}+\frac{1}{R e_{r}}\left[\frac{\partial \tau_{x y}}{\partial x}+\frac{\partial \tau_{y y}}{\partial y}\right.$ $\left.+\frac{\partial \tau_{y z}}{\partial z}\right]$ (along y-axis)
$\frac{\partial\left(\rho u_{z}\right)}{\partial t}+\frac{\partial\left(\rho u_{x} u_{z}\right)}{\partial x}+\frac{\partial\left(\rho u_{y} u_{z}\right)}{\partial y}+\frac{\partial\left(\rho u_{z}^{2}\right)}{\partial z}=-\frac{\partial p}{\partial t}+\frac{1}{R e_{r}}\left[\frac{\partial \tau_{x z}}{\partial x}+\frac{\partial \tau_{y z}}{\partial y}\right.$
$\left.+\frac{\partial \tau_{z z}}{\partial z}\right]$ (along $z$-axis). These are three equations and unknown are more so for unique and smooth solution, add two more constraints 1) continuity of conservation of momentum and 2) conservation of energy equation given as below
$\frac{\partial(\rho)}{\partial t}+\frac{\partial\left(\rho u_{x}\right)}{\partial x}++\frac{\partial\left(\rho u_{y}\right)}{\partial y}+\frac{\partial\left(\rho u_{z}\right)}{\partial z}=0$ i.e. $\quad \nabla .(\rho . U)=0$.
(Continuity equation)
And
$\frac{\partial\left(E_{t}\right)}{\partial t}+\frac{\partial\left(\rho u_{x} E_{t}\right)}{\partial x}++\frac{\partial\left(\rho u_{y} E_{t}\right)}{\partial y}+\frac{\partial\left(\rho u_{z} E_{t}\right)}{\partial z}=-\frac{\partial\left(u_{x} p\right)}{\partial x}-\frac{\partial\left(u_{y} p\right)}{\partial y}-\frac{\partial\left(u_{z} p\right)}{\partial z}$
$-\frac{1}{R_{e_{r}} P_{r}}\left[\frac{\partial q_{x}}{\partial x}+\frac{\partial q_{y}}{\partial y}+\frac{\partial q_{z}}{\partial z}\right]+\frac{1}{R_{r_{r}}}\left[\frac{\partial\left(u_{x} \tau_{x x}+u_{y} \tau_{x y}+u_{z} \tau_{x z}\right)}{\partial x}\right.$ $\left.+\frac{\partial\left(u_{x} \tau_{x y}+u_{y} \tau_{y y}+u_{z} \tau_{y z}\right)}{\partial y}+\frac{\partial\left(u_{x} \tau_{x z}+u_{y} \tau_{y z}+u_{z} \tau_{z z}\right)}{\partial z}\right]$ (Energy Equation)
Let us write the system of first four of these equations into matrix form as

$$
\begin{aligned}
& R_{e_{r}}\left(\left[\begin{array}{llll}
\frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} & \frac{\partial}{\partial t}
\end{array}\right]\left[\begin{array}{cccc}
\rho u_{x} & \rho u_{y} & \rho u_{z} & \rho \\
\rho u_{x}^{2}+p & \rho u_{x} u_{y} & \rho u_{x} u_{z} & \rho u_{x} \\
\rho u_{x} u_{y} & \rho u_{z}^{2}+p & \rho u_{y} u_{z} & \rho u_{y} \\
\rho u_{x} u_{z} & \rho u_{y} u_{z} & \rho u_{z}^{2}+p & \rho u_{z}
\end{array}\right]\right) \\
& =\left(\left[\begin{array}{llll}
\frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} & \frac{\partial}{\partial t}
\end{array}\right]\left[\begin{array}{cccc}
\tau_{x x} & \tau_{x y} & \tau_{x z} & 0 \\
\tau_{x y} & \tau_{y y} & \tau_{y z} & 0 \\
\tau_{x z} & \tau_{y z} & \tau_{z z} & 0 \\
0 & 0 & 0 & 0
\end{array}\right]\right.
\end{aligned}
$$

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Now the momentum matrix of the above equation is

$$
\begin{aligned}
& {\left[\begin{array}{cccc}
\rho u_{x} & \rho u_{y} & \rho u_{z} & \rho \\
\rho u_{x}^{2}+p & \rho u_{x} u_{y} & \rho u_{x} u_{z} & \rho u_{x} \\
\rho u_{x} u_{y} & \rho u_{y}^{2}+p & \rho u_{y} u_{z} & \rho u_{y} \\
\rho u_{x} u_{z} & \rho u_{y} u_{z} & \rho u_{z}^{2}+p & \rho u_{z}
\end{array}\right]=} \\
& \rho \cdot\left[\begin{array}{cccc}
u_{x} & u_{y} & u_{z} & 1 \\
u_{x}^{2} & u_{x} u_{y} & u_{x} u_{z} & u_{x} \\
u_{x} u_{y} & u_{y}^{2} & u_{y} u_{z} & u_{y} \\
u_{x} u_{z} & u_{y} u_{z} & u_{z}^{2} & u_{z}
\end{array}\right]+\left[\begin{array}{cccc}
0 & 0 & 0 & 0 \\
p & 0 & 0 & 0 \\
0 & p & 0 & 0 \\
0 & 0 & p & 0
\end{array}\right]
\end{aligned}
$$

So determinant of velocity matrix is $=\operatorname{det}$

$$
\begin{aligned}
&\left(\rho \cdot\left[\begin{array}{cccc}
u_{x} & u_{y} & u_{z} & 1 \\
u_{x}^{2} & u_{x} u_{y} & u_{x} u_{z} & u_{x} \\
u_{x} u_{y} & u_{y}^{2} & u_{y} u_{z} & u_{y} \\
u_{x} u_{z} & u_{y} u_{z} & u_{z}^{2} & u_{z}
\end{array}\right]\right) \\
&= \rho^{4} \cdot u_{x} \cdot u_{y} \cdot u_{z} \cdot \operatorname{det}\left(\left[\begin{array}{cccc}
u_{x} & u_{y} & u_{z} & 1 \\
u_{x} & u_{y} & u_{z} & 1 \\
u_{x} & u_{y} & u_{z} & 1 \\
u_{x} & u_{y} & u_{z} & 1
\end{array}\right]\right) \\
&=0
\end{aligned}
$$

Here we can conclude that all the points where fluid motion takes placed are critical points. So the turbulence for motion can occur even if very small external force is applied. Also the divergence matrix is applied on the velocity matrix to get acceleration. So there is scope to study entire description motion of fluid i. e. extreme values, change of nature of path using eigen values and trace of the acceleration matrix. Hassian matrix can be used to find extreme values of motion.

## CONCLUSION

The Nevier-Stokes equations are converted into matrix form. Also all points of the motion through which the fluid is passing are critical points. So the turbulence for motion can occur even if very small force is applied.

## References

1. Patil D. R., Bhadane A. P., Poonia M. S., "Solution of non-linear partial differential equations using modified Charpit's method", International Journal of Engineering Science and Computing, ISSN-2321-3361, Vol. 6, Issue -8, (2016), 2651-2653.
2. Ladyzhenskaya O., "The Mathematical Theory of Viscous Incompressible Flows" (2nd ed.), New York: Gordon and Breach, (1969).
3. Patil D. R., Poonia M. S., "Current Updates in Fluid Dynamics with Special Reference to Europe", International Journal of Mathematics and its Applications, ISSN-2347-1557, Vol. 4, Issue 2B,(2016), 123-125.
4. I. N. Sneddon, "Elements of Partial Differential Equations", McGRAW-Hill international Editions, 1957.
5. D. L. Bernstein, "Existence Theorems in Partial differential Equations", Annals of Mathematics Study, No. 23, Princeton University Press, Princeton, New Jersey, 1950.

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